

A solved problem in rotational dynamics

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Abstract

This paper is a redo of an article that first appeared in the *Arizona Journal of Natural Philosophy*, July 1992. What's special about the solution here is that it uses geometric algebra.¹

Given an axially symmetric rigid body rotating about a fixed point under no torques, and given that the angle between the symmetry axis and the rotational velocity vector $\boldsymbol{\omega}$ is α , show that the tangent of the angle γ between $\boldsymbol{\omega}$ and the angular momentum vector \mathbf{L} is given by

$$\tan \gamma = \frac{(I_s - I) \tan \alpha}{I_s + I \tan^2 \alpha} \quad (1)$$

where I_s is the moment of inertia about the symmetry axis and I is the moment of inertia about the axis normal the symmetry axis, which is the z -axis.

SOLUTION:

Let $\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3$ be an orthonormal frame along principal directions of the body, where \mathbf{e}_3 is along the symmetry axis. Let θ be the angle between the symmetry axis and \mathbf{L} , where

$$\mathbf{L} = I\omega_x\mathbf{e}_1 + I\omega_y\mathbf{e}_2 + I_s\omega_z\mathbf{e}_3 \quad (2)$$

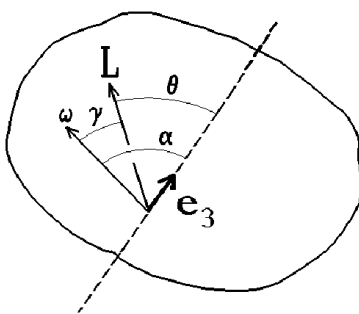


Figure 1. The relevant vectors and angles displayed. The dashed line is the axis of symmetry. Vectors \mathbf{L} , $\boldsymbol{\omega}$, and \mathbf{e}_3 are coplanar.

¹For references to geometric algebra, see the bibliography at the end of the article.

The following results will be useful:

$$\tan \alpha = \frac{|\boldsymbol{\omega} \wedge \mathbf{e}_3|}{\boldsymbol{\omega} \cdot \mathbf{e}_3} = \frac{\sqrt{\omega_x^2 + \omega_y^2}}{\omega_z} \quad (3)$$

$$\tan \theta = \frac{|\mathbf{L} \wedge \mathbf{e}_3|}{\mathbf{L} \cdot \mathbf{e}_3} = \frac{\sqrt{L_x^2 + L_y^2}}{L_z} = \frac{I}{I_s} \tan \alpha \quad (4)$$

$$\tan \gamma = \frac{|\mathbf{L} \wedge \boldsymbol{\omega}|}{\mathbf{L} \cdot \boldsymbol{\omega}} \quad (5)$$

Thus we can solve for $\tan \gamma$ in terms of \mathbf{a} if we can solve for $\mathbf{L} \wedge \boldsymbol{\omega}$ and $\mathbf{L} \cdot \boldsymbol{\omega}$ in terms of α . The following step is key—a virtual emplacement based on the fact that $\mathbf{e}_3^2 = 1$ and that the geometric product is associative.

$$\mathbf{L}\boldsymbol{\omega} = \mathbf{L}(\mathbf{e}_3\mathbf{e}_3)\boldsymbol{\omega} = (\mathbf{L}\mathbf{e}_3)(\mathbf{e}_3\boldsymbol{\omega}) \quad (6)$$

$$\mathbf{L} \cdot \boldsymbol{\omega} + \mathbf{L} \wedge \boldsymbol{\omega} = (\mathbf{L} \cdot \mathbf{e}_3 + \mathbf{L} \wedge \mathbf{e}_3)(\mathbf{e}_3 \cdot \boldsymbol{\omega} + \mathbf{e}_3 \wedge \boldsymbol{\omega}) \quad (7)$$

On equating scalar and bivector parts, respectively, we get

$$\mathbf{L} \cdot \boldsymbol{\omega} = \mathbf{L} \cdot \mathbf{e}_3 \mathbf{e}_3 \cdot \boldsymbol{\omega} + |\mathbf{L} \wedge \mathbf{e}_3| |\boldsymbol{\omega} \wedge \mathbf{e}_3| = L_z \omega_z + |\mathbf{L} \wedge \mathbf{e}_3| |\boldsymbol{\omega} \wedge \mathbf{e}_3|, \quad (8)$$

$$\mathbf{L} \wedge \boldsymbol{\omega} = \mathbf{L} \cdot \mathbf{e}_3 \mathbf{e}_3 \wedge \boldsymbol{\omega} - \boldsymbol{\omega} \cdot \mathbf{e}_3 \mathbf{L} \wedge \mathbf{e}_3 = L_z \mathbf{e}_3 \wedge \boldsymbol{\omega} - \omega_z \mathbf{L} \wedge \mathbf{e}_3. \quad (9)$$

Since all the bivectors of the last equation are coplanar, we can easily scalarize it, we get

$$|\mathbf{L} \wedge \boldsymbol{\omega}| = \mathbf{L} \cdot \mathbf{e}_3 |\mathbf{e}_3 \wedge \boldsymbol{\omega}| - |\boldsymbol{\omega} \cdot \mathbf{e}_3| |\mathbf{L} \wedge \mathbf{e}_3| = L_z |\mathbf{e}_3 \wedge \boldsymbol{\omega}| - \omega_z |\mathbf{L} \wedge \mathbf{e}_3|. \quad (10)$$

So that

$$\begin{aligned} \tan \gamma &= \frac{L_z |\mathbf{e}_3 \wedge \boldsymbol{\omega}| - \omega_z |\mathbf{L} \wedge \mathbf{e}_3|}{L_z \omega_z + |\mathbf{L} \wedge \mathbf{e}_3| |\boldsymbol{\omega} \wedge \mathbf{e}_3|} \\ &= \frac{|\mathbf{e}_3 \wedge \boldsymbol{\omega}|/\omega_z - |\mathbf{L} \wedge \mathbf{e}_3|/L_z}{1 + \frac{|\mathbf{L} \wedge \mathbf{e}_3|}{L_z} \frac{|\boldsymbol{\omega} \wedge \mathbf{e}_3|}{\omega_z}} \\ &= \frac{(I_s - I) \tan \alpha}{I_s + I \tan^2 \alpha} \end{aligned} \quad (11)$$

And this concludes the proof.

David Hestenes and Garret Sobczyk. 1984. *Clifford Algebra to Geometric Calculus*. D. Reidel Publishing Co. Dordrecht, Holland.

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